

Temperature Dependence of Channel Mobility in HfO₂-Gated NMOSFETs

W. J. Zhu, *Member, IEEE*, and T. P. Ma, *Fellow, IEEE*

Abstract—The degradation mechanisms of effective electron channel mobility in HfO₂-gated nMOSFETs have been studied by analyzing experimental data at various temperatures from 120 to 320 K. The major finding is that, while significant Coulomb scattering plays an important role in causing the observed mobility degradation, it does not account for all of the degradation; rather, it requires an extra phonon scattering mechanism, beyond that arising from the phonons in the Si substrate, to explain our experimental results. This extra phonon scattering mechanism has been found to exhibit relatively weak temperature dependence, and is attributed to the soft optical phonons in the HfO₂ layer.

Index Terms—Channel mobility, Coulomb scattering, high- κ dielectrics, MOSFETs, phonon scattering.

I. INTRODUCTION

TO REALIZE the scaling scenario projected in the International Technology Roadmap for Semiconductors (ITRS), it is widely believed that a high- κ (high permittivity) dielectric is needed to serve as the CMOS gate dielectric to reduce significantly the gate leakage current [1]. After extensive research by numerous groups, several high- κ dielectrics have shown promising results [2]–[4]. However, there are still many challenges that have held back the actual implementation of these promising candidates. One of the major challenges is the significantly degraded channel mobility for high- κ gated MOSFETs compared to their SiO₂ counterparts [4], and there is not yet a clear understanding of the causes behind this degraded mobility. Coulomb scattering due to high density of interface trapped charges and fixed oxide charges in high- κ dielectrics appears to be an important contributor [5]. However, our preliminary analyzes of many high- κ samples have suggested that increased Coulomb scattering alone cannot account for all of the mobility degradation, and we believe that additional scattering mechanisms may be responsible for the difference. From a systematic study of the field dependence and temperature dependence of channel mobility in a set of HfO₂-gated nMOSFETs, we obtained strong evidence to support the “remote phonon” scattering mechanism, i.e., scattering by soft optical phonons in the high- κ dielectric, proposed by Fischetti, *et al.* [6]. This letter will report some of the key results that we have obtained through this study.

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II. ACCURATE CHANNEL MOBILITY MEASUREMENT

To study the scattering mechanisms, it is of crucial importance to accurately measure the channel mobility. In this letter, we obtained the effective channel mobility, μ_{eff} , from the drain current, I_d , in the linear region from the following equation [7]:

$$\mu_{\text{eff}} = \frac{L}{W} \cdot \frac{I_d(V_g)}{V_d Q_{\text{inv}}(V_g)} \quad (1)$$

where the inversion charge density, Q_{inv} , is usually extracted by measuring the gate-channel capacitance C_{gc} as a function of gate voltage, V_g , by the split-capacitance-voltage (C - V) technique [8], and integrating

$$Q_{\text{inv}} = \int_{-\infty}^{V_g} C_{\text{gc}}(V_g) dV_g. \quad (2)$$

It has been shown that, because of significant trapping of carriers in high- κ gated MOSFETs, the use of the conventional split- C - V technique [9] will result in significant overestimate of Q_{inv} , and thus a significant underestimate of the effective mobility [9]. Therefore, in this study we used a modified split- C - V technique [9] to obtain the Q_{inv} value accurately, which allowed us to extract the effective channel mobility accurately even in the presence of significant carrier trapping.

III. DEVICE FABRICATION AND MEASUREMENTS

Nonself-aligned nMOSFETs are fabricated with ultrathin [with equivalent oxide thickness (EOT) ~ 2 nm] HfO₂ gate dielectrics on p-Si substrates with a dopant concentration of approximately $4 \times 10^{15}/\text{cm}^3$. The source and drain regions of the nMOSFETs are implanted with phosphorous. The implantation dose is $5 \times 10^{15}/\text{cm}^2$, and the implantation energy is 80 keV. After source/drain activation at 1000 °C, thin HfO₂ films are deposited by jet-vapor deposition (JVD) [10], [11] at room temperature on HF-last Si substrates. The postdeposition anneal (PDA) is done in forming gas at 700 °C. Aluminum is used as the gate electrode and source/drain contacts. The split- C - V and I_d - V_g characteristics are measured at 320 to 120 K in a vacuum of 20 mTorr.

IV. RESULTS AND DISCUSSION

Fig. 1 displays the effective mobility, μ_{eff} , of a HfO₂-gated nMOSFET measured at various temperatures from 120 to 320 K. For comparison, the universal mobility for SiO₂-gated nMOSFET is also shown in the same temperature range. In Fig. 1 the lowest set of curves, marked “HfO₂ w/o correction”, was obtained by the use of the conventional split- C - V technique to get Q_{inv} from (2), which underestimates the effective mobility because of trapping. By the use of the modified split- C - V

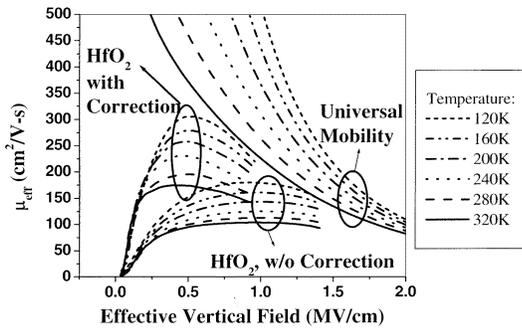


Fig. 1. Effective electron mobility in HfO₂-gated nMOSFET at various temperatures from 120 to 320 K, with and without interface-trap correction, as compared to universal mobility over the same temperature range.

technique to measure the Q_{inv} correctly, we obtained the set of curves marked “HfO₂ with correction” in Fig. 1, which is substantially higher than the uncorrected set, but is still much lower than that for the SiO₂-gated universal curves.

Fig. 2 shows the effective mobility as a function of effective field at 300 K for an nMOSFET with HfO₂ as gate dielectric and another one with SiO₂ as gate dielectric, both with Al gate. It is apparent that, at low fields the effective mobility of the HfO₂ sample is significantly lower than that of the SiO₂ sample, which may be attributed to the much stronger Coulomb scattering for the former due to its much higher trapped charge density (include interface traps, border traps [12] and bulk traps), as well as fixed oxide charge density. Here the Coulomb scattering includes possible remote coulomb scattering RCS [13] due to charges at some distances away from the dielectric/substrate interface. Note that in the low-field regime, the effective mobilities for these samples are approximately linearly proportional to the inversion charge density, which is consistent with the carrier screening effect [14].

In addition to Coulomb scattering caused by high densities of charges, the scattering due to soft optical phonons in high- κ dielectrics [6] is an intriguing possibility that cannot be overlooked. To analyze the possible phonon scattering effect, we make use of the following equation, according to Matthiessen’s rule:

$$\frac{1}{\mu_{ph}} = \frac{1}{\mu_{eff}} - \frac{1}{\mu_{coul}} - \frac{1}{\mu_{sr}} \quad (3)$$

where μ_{eff} is the total effective mobility μ_{coul} is the mobility limited by Coulomb scattering, and μ_{sr} is the mobility limited by surface roughness. To obtain the values for μ_{ph} , we extracted the values for μ_{eff} from split- $C-V$ with interface-trap correction, those for μ_{coul} from the effective mobility at low fields by linear fitting [14] and then extrapolated to high fields, and μ_{sr} by using the following equation [14]:

$$\mu_{sr} = B \cdot E_{eff}^{-2.6} \quad (4)$$

where B is 4.5×10^{19} . The unit for E_{eff} is V/m, and the unit for μ_{sr} is m²/V - s. Here we assume that the surface roughness for HfO₂ is similar to SiO₂, based on the fact that high-resolution transmission electron microscope (TEM) images of our HfO₂/Si samples show very smooth dielectrics/Si interface (data not shown), and the atomic force microscopy (AFM) images of our HfO₂ sample showed root-mean-square (rms) of

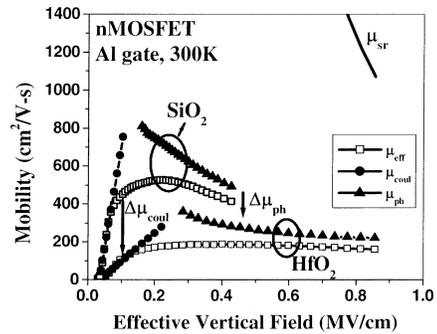


Fig. 2. Effective electron mobility, μ_{eff} , the mobility limited by Coulomb scattering, μ_{coul} , the mobility limited by phonon scattering, μ_{ph} , and that limited by surface roughness, μ_{sr} , for a HfO₂ sample and a SiO₂ sample, where the effective mobility, μ_{eff} , is extracted from split- $C-V$ with interface-trap correction, μ_{coul} is extracted from the effective mobility at low fields by linear fitting, μ_{sr} is calculated according to (4), and μ_{ph} is extracted according to (3). Note that μ_{sr} is much higher than μ_{coul} and μ_{ph} in the field range of interest, and therefore μ_{sr} does not affect μ_{eff} significantly in this range.

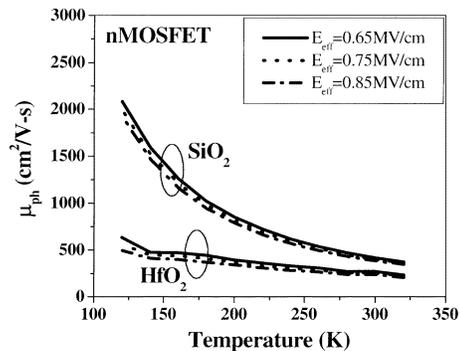


Fig. 3. Mobility limited by phonon scattering, μ_{ph} , for nMOSFET with HfO₂ or SiO₂ as gate dielectric at various effective electric fields from 0.65 mV/cm to 0.85 MV/cm.

0.12 nm, suggesting that the surface roughness of the HfO₂ sample is comparable to its SiO₂ counterpart. In any case, in the mid-to-high field range we studied, the surface roughness scattering is much smaller than phonon scattering and Coulomb scattering, so any reasonable difference in μ_{sr} will not impact significantly the extracted μ_{ph} according to Matthiessen’s rule.

Fig. 2 shows the extracted mobility limited by phonon scattering for a HfO₂-gated MOSFET and that for its SiO₂ counterpart. As we can see, the mobility limited by phonon scattering for the HfO₂-gated MOSFET is significantly lower than that for its SiO₂ counterpart, indicating a more severe phonon scattering for the former. Since the phonon scattering from the silicon substrate should be the same for both samples, this result suggests that the HfO₂-gated MOSFET has an additional source of phonon scattering, which is consistent with the high- κ related soft optical phonon model proposed by Fischetti [6].

The temperature dependence of the extracted mobility limited by phonon scattering, μ_{ph} , for HfO₂ and SiO₂ at various fields from 0.65 MV/cm to 0.85 MV/cm is shown in Fig. 3. Here the SiO₂ curves are calculated by using the equation: $\mu_{ph} = AE_{eff}^{-0.3}T^{-1.75}$, where A is 2×10^5 in MKS unit [14]. As we can see, the mobility limited by phonon scattering for HfO₂ is much lower than that for SiO₂ for all temperatures we tested here, and the lower the temperature, the larger the degradation of μ_{ph} for HfO₂ as compared to SiO₂.

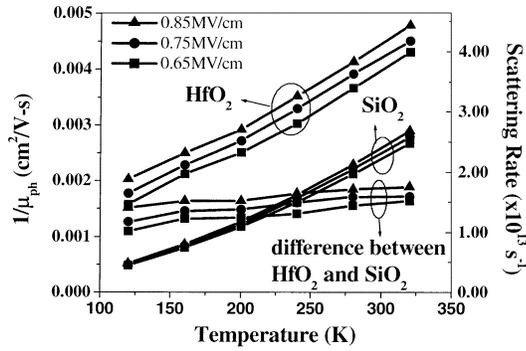


Fig. 4. Scattering rate due to phonons for nMOSFET with HfO₂ or SiO₂ as gate dielectric, indicating additional scattering in the HfO₂ sample as compare to the SiO₂ sample. The difference between the two is also plotted, and is attributed to scattering by soft optical phonons in HfO₂.

To investigate what may have caused the difference, we plot in Fig. 4 the reciprocal data of Fig. 3, or the phonon scattering rate as a function of temperature. Here the scattering rate is calculated according to the equation: $\text{Scattering_Rate} = (q/(\mu \cdot m^*))$, where m^* is the effective mass of the carrier in the Si inversion channel, and is taken to be $m^* = 0.91m_0$, where m_0 is the free electron mass. One can see that the phonon scattering rate for the HfO₂-gated sample is much higher than that for its SiO₂-gate counterpart. Since both the HfO₂-gated and the SiO₂-gated MOSFETs have essentially identical Si substrates, we attribute the difference in phonon scattering rates to other phonon sources, and the most likely source is the soft optical phonons in the high- κ gate dielectric, as proposed by Fischetti *et al.* [6]. By taking the difference between the two sets of curves in Fig. 4, one can get the scattering rate caused by this additional phonon source in the HfO₂ sample, and the resulting curves are also plotted in Fig. 4.

As one can see in Fig. 4, the scattering rate by the soft optical phonons in HfO₂ is a relatively weak function of temperature. The weak temperature dependence of the scattering rate for soft optical phonon can be qualitatively explained by the following derivation. In the two dimensional deformation potential theory of surface phonon scattering, the scattering rate due to optical phonon may be expressed as [15]

$$\frac{1}{\tau_{op}} \propto [N_R + (N_R + 1)u(E - \hbar\omega)] \quad (5)$$

where $N_R = (1/(e^{\hbar\omega/kT} - 1))$ is the phonon occupation number, E is the carrier energy, ω is the phonon frequency, and $u(x)$ is the unit step function, $u(x < 0) = 0$ and $u(x > 0) = 1$. Assume that the phonon energy is smaller than the thermal carrier energy (see discussion below), i.e., $\hbar\omega < E$, then

$$\frac{1}{\tau_{op}} \propto (2N_R + 1) = \frac{e^x + 1}{e^x - 1} \quad \text{where } x = \frac{\hbar\omega}{kT}. \quad (6)$$

Equation (6) indicates that when $\hbar\omega \ll kT$, $(1/\tau_{op}) \propto T$, and when $\hbar\omega > kT$, $(1/\tau_{op})$ approaches a constant, i.e independent of temperature. For HfO₂, the soft optical phonon energies are 12.40 meV (TO1) and 48.35 meV (TO2) [6], which are to be compared to the kT range of 10.4 meV to 27.6 meV for our measurement temperatures (120 K to 320 K). Therefore, one would expect a temperature independent scattering rate for the TO2 phonons, and a constant-to-linear temperature dependence for the TO1 phonons, which is consistent with the weak temperature dependence shown in Fig. 4.

V. SUMMARY

The degradation mechanisms of effective channel mobility in HfO₂-gated NMOSFETs have been studied at various temperatures from 120 to 320 K. We have found that, while significant Coulomb scattering does bring down the effective channel mobility substantially in high- κ gated NMOSFET, it cannot account for all of the mobility degradation as compared to the universal mobility curve. A systematic study of the temperature dependence and field dependence of the effective mobility has led us to conclude that an additional scattering mechanism, that due to the soft optical phonons in the HfO₂ gate dielectric, may be responsible for the enhanced mobility degradation in HfO₂-gated MOSFETs. The scattering rate due to the soft optical phonons in HfO₂ has been found to be relatively weakly dependent on temperature, which is consistent with our simple theoretical calculation.

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